

# Observation, Characterization, and Mitigation of the Internal $p$ - $n$ Junction in Pyrite $\text{FeS}_2$ , a Potential Low-cost Solar Absorber

**B. Voigt,<sup>1+</sup> W. Moore,<sup>1</sup> J. Walter,<sup>1</sup> B. Das,<sup>1</sup> M. Maiti,<sup>1</sup> M. Manno,<sup>1</sup>  
E. S. Aydil,<sup>1,2</sup> C. Leighton<sup>1</sup>**

<sup>1</sup> Department of Chemical Engineering and Materials Science, University of Minnesota,  
421 Washington Ave SE, Minneapolis, MN 55455, USA

<sup>2</sup> Department of Chemical and Biomolecular Engineering, New York University Tandon  
School of Engineering, 6 MetroTech Center, Brooklyn, NY 11201, USA

Pyrite  $\text{FeS}_2$  is widely acknowledged as an ideal semiconductor for thin film solar cells due to its earth-abundance, low toxicity, low cost, suitable band gap (0.95 eV) and minority carrier diffusion length, and high visible light absorptivity. Power conversion efficiencies of  $\text{FeS}_2$  heterojunction solar cells, however, have never exceeded 3% due to low open-circuit voltages ( $V_{\text{OC}} < 0.3$  V). One hypothesis emerging from recent temperature ( $T$ )-dependent transport measurements of high quality single crystals is that this low  $V_{\text{OC}}$  is due to a conductive pyrite surface with a carrier type ( $p$ -type) inverted from bulk ( $n$ -type) [1,2]. This could create a leaky (*i.e.*, low- $V_{\text{OC}}$ ) internal  $p$ - $n$  junction, thus limiting heterojunction solar cell efficiencies. These studies established conduction through a 1-3 nm-thick,  $p$ -type surface upon freeze-out of  $n$ -type bulk carriers [1,2]. Two parallel resistors representing surface and bulk conduction can describe the  $T$ -dependence of resistivity across a wide  $T$  range (50-500 K) [1] and the non-linear Hall effect observed near the crossover between bulk- and surface-dominated conduction upon cooling below 300 K [2]. Notably, what has neither been observed nor characterized, however, is the internal  $p$ - $n$  junction implied by this  $p$ -type surface and  $n$ -type bulk. Here, we directly observe this internal junction for the first time. In-plane sheet resistance ( $R_S$ ) measurements of polished crystals doped heavily  $n$ -type *via* sulfur vacancies are shown to display an effect where metallic-like transport abruptly transitions to rapidly increasing  $R_S$  below  $\sim 175$  K, eventually transitioning to surface conduction at lower  $T$  ( $< 100$  K). We show that this very unusual  $T$ -dependence can be well described by incorporating an exponentially- $T$ -dependent junction resistance into the parallel resistor model. Junction barrier heights extracted from the model are typically 0.15 – 0.30 eV, in good agreement with typical  $V_{\text{OC}}$  values in past heterojunction solar cells, suggesting that this internal junction may, in fact, be limiting conversion efficiencies. Interestingly, while junction influence in  $R_S(T)$  is independent of contact materials such as In, Ag, Fe, Co, and Ni,  $\text{CoS}_2$  contacts mitigate this junction, allowing the first characterization of bulk properties to low  $T$ . Access to bulk properties at low  $T$  unveils rich phenomena, such as the onset of a smaller donor activation energy below 175 K, non-linear Hall effect near 100 K, and an unusual resistivity anomaly at  $T \leq 10$  K, showcasing  $\text{CoS}_2$  contacts as a way to both mitigate this junction and advance understanding of electronic transport in  $\text{FeS}_2$ . This work was supported by the customers of Xcel Energy through a grant from the Renewables Development Fund.

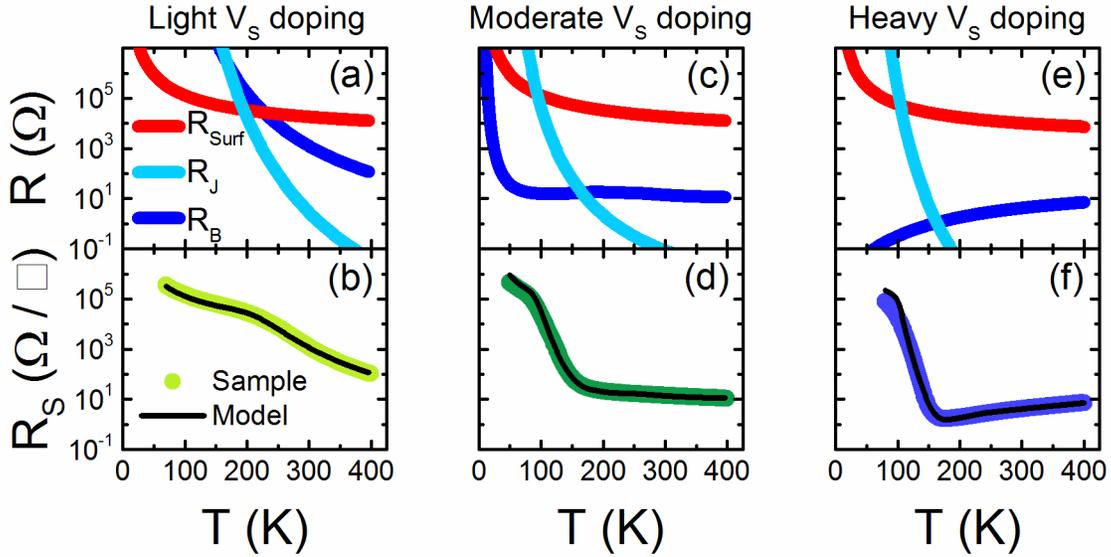
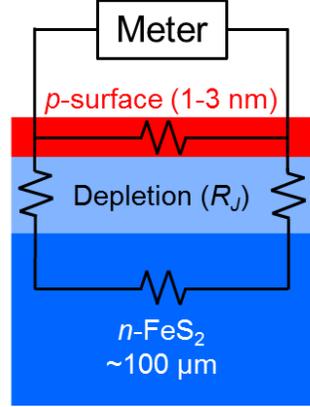
[1] M. Limpinsel *et al.*, Energy Environ. Sci. **7**, 1974 (2014).

[2] J. Walter, *et al.*, Phys. Rev. Mater. **1**, 065403 (2017).

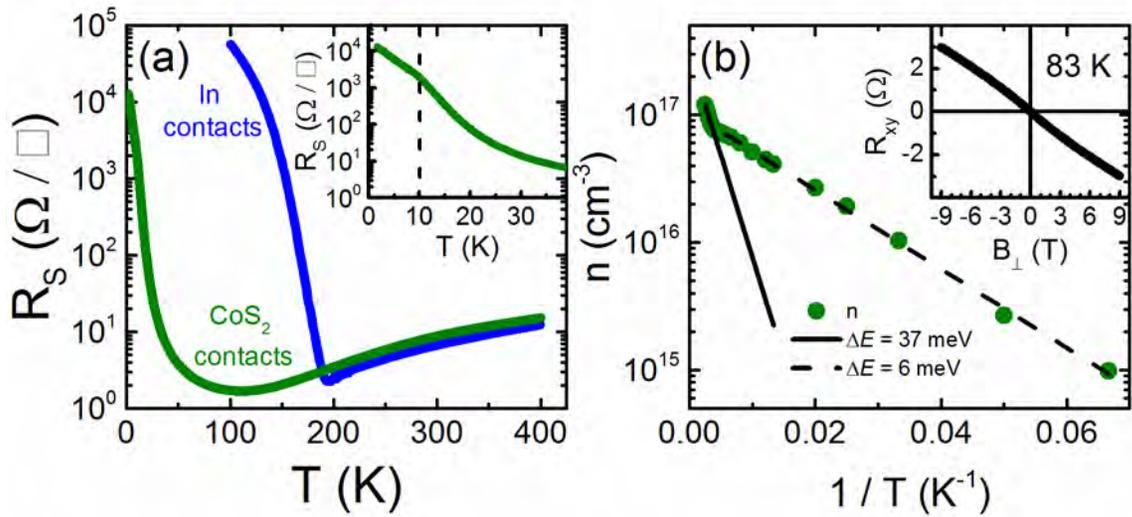
<sup>+</sup> Author for correspondence: voigt132@umn.edu

## Supplementary Pages

**Figure 1.** Expected resistor network in horizontal transport measurements of a pyrite  $\text{FeS}_2$  single crystal ( $\sim 100 \mu\text{m}$  thick) with an  $n$ -type interior (shown in blue) and a degenerately-doped,  $p$ -type surface (shown in red). The depletion region on the  $n$ -type side of the internal junction, and the resistance associated with it ( $R_J$ ), is shown in light blue.



**Figure 2.** (a,c,d) The temperature ( $T$ ) dependence of surface ( $R_{\text{Surf}}$ ), junction ( $R_J$ ), and bulk resistance ( $R_B$ ) contributions to the total sheet resistance ( $R_S$ ) of  $\text{FeS}_2$  crystals grown under light (a), moderate (c), and heavy (e) sulfur vacancy ( $V_S$ ) doping. The  $T$ -dependence of  $R_{\text{Surf}}$  (red data) is independent of  $V_S$  doping.  $R_B(T)$  (blue data) is quantitatively described in each case using the Hall effect and the Drude model, by which the  $T$ -dependences of electron density and mobility are described by simple activated and power law behavior, respectively (data not shown).  $R_J(T)$  (light blue data) is described assuming a Schottky junction, with  $R_J(T) = R_{0,J} e^{q\phi_B/k_B T}$ , where  $R_{0,J}$  is a pre-exponential factor,  $q$  is the electric charge,  $k_B$  is Boltzmann's constant, and  $\phi_B$  is the Schottky barrier height (in eV). (b,d,f) Comparison of the  $T$ -dependence of measured sheet resistance ( $R_S$ , colored data) and calculated sheet resistance (black lines). In each case, the measured  $R_S$  is well described by the parallel resistor network shown in Figure 1, adjusted from previous work by including  $R_J$ . Interestingly,  $R_J$  is only made manifest in crystals doped with larger  $V_S$  concentrations, where  $R_B$  does not freeze out upon cooling to intermediate  $T$  ( $\sim 200$  K) due to donor band broadening and an evolution towards an insulator-metal transition.



**Figure 3.** (a)  $R_S(T)$  of an  $\text{FeS}_2$  crystal heavily doped with  $\text{V}_S$ , contacted using In soldered (blue data) vs.  $\text{CoS}_2$  (green data) contacts.  $\text{CoS}_2$  contacts mitigate the strong  $T$ -dependent  $R_J$ , allowing access to bulk  $\text{FeS}_2$  transport at lower  $T$ . Shown in the inset is bulk  $\text{FeS}_2$  transport at low  $T$ , accessed using  $\text{CoS}_2$  contacts, which highlights the resistivity anomaly near and below 10 K (vertical dashed line).  $\text{CoS}_2$  is a ferromagnetic metal, also crystallizing in the pyrite structure, and Co is a known shallow donor in  $\text{FeS}_2$ . The current hypothesis is that, through the mild heat treatment (350 °C, 8 hrs, in S vapor) we use to sulfidize sputtered Co contacts into  $\text{CoS}_2$ , Co in-diffuses, strongly doping the near-surface region of the  $\text{FeS}_2$  crystal and shorting the junction resistance. (b) Arrhenius plot of electron density ( $n$ ) vs.  $1/T$ , where the slope is proportional to the activation energy ( $\Delta E$ ) of the donor state contributing to extrinsic conduction. A smaller  $\Delta E$  (6 meV) is observed below  $\sim 175$  K; this is not accessible without  $\text{CoS}_2$  contacts. Inset: the non-linear magnetic field dependence of the Hall resistance ( $R_{xy}$ ) at 83 K.