

## Advanced Characterization Techniques

Room Davis Hall 101 - Session AC+TM-MoM

### Characterization/Modeling I

Moderators: Michael Scarpulla, University of Utah, Uttam Singiseti, University of Buffalo, SUNY

9:15am AC+TM-MoM-4 Electric Field Induced Defect Redistribution at Ni-Ga<sub>2</sub>O<sub>3</sub> Interfaces, *Daram Ramdin, H. Huang, S. Dhara, S. Rajan, J. Hwang, L. Brillson*, The Ohio State University

$\beta$ -Ga<sub>2</sub>O<sub>3</sub> is a prime semiconductor for high power electronics due to its high intrinsic critical field of  $\sim 8$  MV/cm. However, premature dielectric breakdown that occurs at lower field gradients remains a challenge that is relatively unexplored at the atomic and nanoscale. Here we use UHV depth-resolved cathodoluminescence (DRCLS) correlated with scanning transmission electron microscopy (STEM) to describe how high electric fields in a Ni/Au Schottky diode on HVPE (001) Ga<sub>2</sub>O<sub>3</sub> cause atomic lattice arrangements, depth-dependent phase changes, Ni diffusion and native defect rearrangements all on a nanometer scale and a function of increasing electric field gradient up to and past dielectric breakdown.

Before electrical fields are applied, DRCLS near the Ni/  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> interfacial region exhibits new, above  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> band gap 5.29 eV and 5.82 eV CL emissions corresponding to a  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> phase. Also present are a set of characteristic intrinsic  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> defect emissions including a  $\sim 2.4$  eV feature which increases under the contact. STEM confirms a defective misaligned  $\gamma$  phase that is of average thickness  $\sim 5$  nm.

With increasing reverse bias at 143 nm below the M-S interface, in-situ DRCLS shows increasing above-bandgap emissions relative to the intrinsic 3.2 eV and 3.6 eV emissions. Prior to breakdown, these  $\sim 5.2$  eV and  $\sim 5.8$  eV emissions remain without applied fields, indicating Ni diffusion and accompanying local phase inhomogeneities. After breakdown, these emissions are reduced, indicating reduced Ni and  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> phase present at this depth. However, at the intimate interface, STEM imaging confirms that the defective layer thickness triples, with little Ni diffusion observed outside of the defective layer after breakdown.<sup>1</sup> These observations are consistent with Ni diffusion into  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> with metallization and subsequent reverse bias, then diffusion back towards the defective layer during the breakdown process, similar to the effect of annealing on this defective layer (unpublished). Breakdown experiments performed in air show that  $V_{br}$  is 2.7x higher, indicating that thermal effects play a more pronounced role in the breakdown process in UHV. These results provide evidence that dielectric breakdown at the widely used Ni/  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> contact involves an interplay of nanoscale phase change, Ni diffusion, and defect rearrangement.

Support acknowledged from AFOSR Grant No. FA9550-18-1-0066 (DNR & LJB) and AFOSR (GAME MURI) Grant No. FA9550-18-1-0479 (HLH, JH).

1. J. Shi et al., Appl. Mater. Int. 2021, 13, 29083-29092.

9:30am AC+TM-MoM-5 Charge State Transition Levels of Ni in  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> Crystals from Experiment and Theory: Eminently Suitable Candidate for Compensation, *Palvan Seyidov*, Leibniz-Institut für Kristallzüchtung, Germany; *J. Basile Varley*, Lawrence Livermore National Laboratory; *Z. Galazka, T. Chou, A. Popp, K. Irmischer, A. Fiedler*, Leibniz-Institut für Kristallzüchtung, Germany

$\beta$ -Ga<sub>2</sub>O<sub>3</sub> has emerged as a next-generation high-power application due to its large bandgap of 4.85eV and a high theoretical breakdown field of  $\sim 8$  MV/cm, which already resulted in established power rectifiers and MOSFETs with excellent characteristics.<sup>1</sup> Bulk single crystals for substrates can be grown with EFG<sup>2</sup> and Czochralski<sup>3</sup> methods. One of the critical criteria for achieving excellent characteristic in lateral devices is the choice of high-resistive free-standing substrates. In this respect, deep acceptor dopants are necessary to counter the unintentional donor impurities. Recently, Mg, Fe, and Co dopants have been proposed as potential deep acceptors for producing semi-insulating  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals. However, if the acceptor level positions above the mid-band gap, it can lead to electron conduction and loss of the semi-insulating state at high fields or high temperatures, as is the case for Fe and Co dopants. On the other hand, the holes are not mobile (in VB) in  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals resulting in a high semi-insulating state if the acceptor level positions below the mid-band gap, as is the case in the Mg acceptor dopant. However, the acceptor level of Mg positions only  $\sim 1.2$  eV above the VBM, where thermal ionization is to be expected at higher temperatures leading to loss of the semi-insulating state of substrates. Thus, an acceptor dopant is needed that has its acceptor level

below and very close to the mid-band gap, which is the key factor for producing reliable semi-insulating  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals.

We measured the photoconductivity, optical absorption, and temperature-dependent resistivity (up to 1100K) of Ni-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystal grown by the Czochralski method<sup>3</sup>. Fitting results of photoconductivity measurement allow us to identify the energy of the  $E_{ZPL}$ ,  $D_{FC}$ , and  $hw$  of Ni-related deep levels. The first-principal calculations based on DFT support our identifications. Spectral regions observed in optical absorption mainly arise due to the charge transfer from Ni-related deep levels and CB or VB. High-temperature resistivity shows a thermal activation energy of  $\sim 2.0$  eV. Conclusively, from the experiment and theory, a consistent energy scheme: an acceptor level of  $\sim 1.9$  eV (above the VBM), and a donor level of  $\sim 3.8$  eV (below the CBM) were identified. Due to the position of the acceptor level (below and close to the mid-band gap), Ni seems to be a promising acceptor dopant for producing semi-insulating  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> substrates for lateral power devices.

<sup>1</sup> A. J. Green et al., APL Mater **10**(2), 029201 (2022). <sup>2</sup> A. Kuramata et al., Jpn. J. Appl. Phys. **55**(100), 1202A2 (2016). <sup>3</sup> Z. Galazka; J. Appl. Phys. **131** (2022) 031103.

9:45am AC+TM-MoM-6 Comparative Study of Temperature-Dependent Bandgap Transitions in Ga<sub>2</sub>O<sub>3</sub> Polymorphs, *Benjamin M. Janzen, N. Hajizadeh, M. Meißner, M. Marggraf, C. Hartung*, Technical University of Berlin, Germany; *Z. Galazka*, Leibniz-Institut für Kristallzüchtung, Berlin, Germany; *P. Mazzolini, A. Sacchi, R. Fornari*, Department of Mathematical, Physical and Computer Sciences, University of Parma, Italy; *C. Petersen, H. von Wenckstern, M. Grundmann*, Universität Leipzig, Felix-Bloch-Institut für Festkörperphysik, Germany; *E. Kluth, M. Feneberg, R. Goldhahn*, Otto-von-Guericke-University Magdeburg, Germany; *T. Oshima*, Department of Electrical and Electronic Engineering, Saga University, Japan; *T. Kato, H. Nishinaka*, Faculty of Electrical Engineering and Electronics, Kyoto Institute of Technology, Japan; *J. Varley*, Lawrence Livermore National Laboratory; *M. Wagner*, Paul-Drude-Institut für Festkörperelektronik, Germany

The temperature dependence of the optical bandgap has rarely been investigated experimentally for the different polymorphs of Ga<sub>2</sub>O<sub>3</sub>. A direct comparison of the temperature dependence as well as the electron-phonon coupling strengths is made considerably more difficult by the different experimental methods (e.g., reflection spectroscopy, absorption spectroscopy or ellipsometry) used to study the various polymorphs. In particular, there is no study in the literature that provides a self-consistent comparison between the band gap values and the electron-phonon coupling strengths of the different polymorphs using the same experimental technique.

We provide a combined experimental-theoretical study to investigate the electronic bandgap transitions in monoclinic [1]  $\beta$ -, orthorhombic rotational-domain [2] as well as single-domain [3]  $\kappa$ -, rhombohedral [4]  $\alpha$ -, defective spinel [5]  $\gamma$ - and cubic bixbyite [6]  $\delta$ -Ga<sub>2</sub>O<sub>3</sub> as a function of the sample temperature. Temperature-dependent UV photoluminescence excitation (PLE) spectroscopy is employed in the temperature range between 5 K and 300 K and the obtained bandgap energies are compared with room temperature measurements of the dielectric function as determined by spectroscopic ellipsometry. The temperature dependencies are discussed in conjunction with DFT calculations regarding the effects of electron-phonon coupling and the averaged phonon energies.

At 5K, we find that  $\gamma$  and  $\alpha$  possess the largest bandgap energy values around 5.36 eV, with the monoclinic  $\beta$ -polymorph's observed polarization-dependent direct band-to-band transitions exhibiting the smallest bandgap energies between 4.72 eV and 4.99 eV. Regarding the strength of the electron-phonon coupling, we observe the strongest coupling for  $\gamma$  or weakest coupling for  $\kappa$  and  $\delta$ , whereas the interaction appears similarly intense for  $\alpha$  with respect to  $\beta$ . In contrast to the rotational domain structured  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> thin film, the single-domain film reveals a directional dependence of the energy bandgap when polarizing the incident light along the crystallographic a- or b-directions, respectively.

[1]: Z. Galazka, Semicond. Sci. Technol., **33**(11), 113001 (2018).

[2]: P. Mazzolini, B. M. Janzen et al., Adv. Funct. Mater., **33**(2), 2207821 (2023).

[3]: H. Nishinaka et al., Jpn. J. Appl. Phys., **61**(1), 018002 (2022).

[4]: S. Vogt et al., Phys. Status Solidi A, **220**(3), 2200721 (2023).

[5]: L. E. Ratcliff, B. M. Janzen et al., Adv. Mater., **34**(37), 2204217 (2022).

[6]: T. Kato et al., ACS Appl. Electron. Mater., **5**(3), 1715 (2023).

# Monday Morning, August 14, 2023

10:00am **AC+TM-MoM-7 Strain and Composition Dependencies in  $(\text{Al}_x\text{Ga}_{1-x})_2\text{O}_3$  Alloys**, Rafal Korlacki, J. Knudtson, M. Stokey, M. Hilfiker, University of Nebraska-Lincoln; V. Darakchieva, Lund University, Sweden; M. Schubert, University of Nebraska-Lincoln

Strain caused by the lattice mismatch in heteroepitaxial thin-films can be used to optimize the optical performance, as it has been demonstrated for the ternary system of  $(\text{Al},\text{Ga})\text{N}$ . [1-2] In order to apply the same principle to  $(\text{Al}_x\text{Ga}_{1-x})_2\text{O}_3$  alloys, the details of the strain relationships for various electronic and optical properties as a function of composition are needed. We use symmetry-based analysis on how the energy eigenvalues and other properties depend on the components of stress and strain tensors. [3] Then, we perform density functional theory (DFT) calculations for a representative set of structures realizing different model deformation scenarios for both,  $\text{Ga}_2\text{O}_3$  and  $\text{Al}_2\text{O}_3$ , in monoclinic and rhombohedral phases. We obtain the linear deformation potentials for energy eigenvalues [3,4] and other material properties that can be extracted from first principles calculations, such as band-to-band transitions, effective mid-band energies, refractive indices, components of the dielectric tensors, and effective mass parameters. Then, Vegard's rule allows us to construct a simple universal model of strain and composition dependencies of these properties in heterostructures under specific strain patterns, [5,6] and thus allowing rational strain engineering to aid design of new  $(\text{Al}_x\text{Ga}_{1-x})_2\text{O}_3$ -based electronic and photonic devices.

[1] D. Li, K. Jiang, X. Sun, and C. Guo, *Adv. Opt. Photon.* **10**, 43-110 (2018)

[2] J.-M. Wagner and F. Bechstedt, *Phys. Rev.* **B66**, 115202 (2002)

[3] R. Korlacki, M. Stokey, A. Mock, S. Knight, A. Papamichail, V. Darakchieva, and M. Schubert, *Phys. Rev.* **B102**, 180101(R) (2020)

[4] R. Korlacki, J. Knudtson, M. Stokey, M. Hilfiker, V. Darakchieva, and M. Schubert, *Appl. Phys. Lett.* **120**, 042103 (2022)

[5] R. Korlacki, M. Hilfiker, J. Knudtson, M. Stokey, U. Kilic, A. Mauze, Y. Zhang, J. Speck, V. Darakchieva, and M. Schubert, *Phys. Rev. Appl.* **18**, 064019 (2022)

[6] M. Stokey, R. Korlacki, J. Knudtson, A. Mock, M. Hilfiker, A. Mauze, Y. Zhang, J. Speck, A. Papamichail, S. Knight, V. Darakchieva, and M. Schubert, "Phonon modes and strain effects in  $\beta$ - $(\text{Al}_x\text{Ga}_{1-x})_2\text{O}_3$ ," in preparation

10:15am **AC+TM-MoM-8 10 kV  $\text{Ga}_2\text{O}_3$  Schottky Rectifier Operational at 200 °C**, Yuan Qin, M. Xiao, M. Potter, Y. Ma, Center of Power Electronics Systems, Virginia Polytechnic Institute and State University; J. Spencer, Naval Research Laboratory; Z. Du, Ming Hsieh Department of Electrical Engineering, University of Southern California; A. Jacobs, Naval Research Laboratory; K. Sasaki, Novel Crystal Technology Inc., Japan; H. Wang, Ming Hsieh Department of Electrical Engineering, University of Southern California; M. Tadjer, Naval Research Laboratory; Y. Zhang, Center of Power Electronics Systems, Virginia Polytechnic Institute and State University

This work demonstrates a novel lateral  $\text{Ga}_2\text{O}_3$  Schottky barrier diode (SBD) with a  $BV$  over 10kV and a thermally-stable 10kV blocking at high temperatures up to 200°C. The device design to enable such performance is a NiO-based reduced-surface-field (RESURF) structure that achieves a charge balance with the  $\text{Ga}_2\text{O}_3$  channel at high reverse bias.

Fig. 1(a) and (b) show the schematic and top-view scanning electron microscopy (SEM) image of the RESURF  $\text{Ga}_2\text{O}_3$  SBD, respectively. A gap between the p-NiO layer and cathode ( $L_{pc}$ ) is designed to prevent the possible leakage current and punch-through in NiO.

Fig. 1(c) shows the depth profile of the net donor concentration in the  $\text{Ga}_2\text{O}_3$  epi layer. A total charge density ( $\sigma_n$ ) of  $3.8 \times 10^{12} \text{cm}^{-2}$  is estimated in n- $\text{Ga}_2\text{O}_3$ . A NiO/n<sup>+</sup>- $\text{Ga}_2\text{O}_3$  p-n diode is fabricated to extract the acceptor concentration in NiO, which is revealed to be  $8 \times 10^{17} \text{cm}^{-3}$  at 25°C and shows little change at 200°C (Fig. 1(d)). If a charge imbalance margin below 15% is kept for practical device fabrication, the NiO thickness ( $t_{\text{NiO}}$ ) range is estimated to be 61-82 and 58-78nm for the anode-to-cathode length ( $L_{AC}$ ) of 30 and 50 $\mu\text{m}$ .

Fig. 2(a)-(c) show the simulated electric field (E-field) contours of RESURF SBDs with different  $t_{\text{NiO}}$ . In charge balance condition ( $t_{\text{NiO}}=75\text{nm}$ ), the E-field are more evenly distributed in  $\text{Ga}_2\text{O}_3$  channel. Fig. 2(d) manifests that the NiO RESURF region could increase the device on-resistance ( $R_{\text{on}}$ ) due to the vertical depletion effect.

Fig. 3(a) and (b) show the reverse I-V characteristics of the  $\text{Ga}_2\text{O}_3$  RESURF SBDs with various  $t_{\text{NiO}}$  for  $L_{AC}$  of 30 and 50 $\mu\text{m}$ . The breakdown voltage ( $BV$ ) increases with the increasing  $t_{\text{NiO}}$ , reaching >10kV (the measurement limit of our test setup) at  $t_{\text{NiO}}=75\text{nm}$ , and starts to decrease at larger  $t_{\text{NiO}}$ . The two

sets of devices can be swept repeatedly up to 10kV at 200°C (Fig. 3(c)). Fig. 3(d) suggests the determining impact of charge balance on the  $BV$ .

Fig. 4(a) shows the forward I-V characteristics of the non-RESURF and 75-nm-RESURF SBDs, both with  $L_{AC}=30$  and 50 $\mu\text{m}$ . The larger  $R_{\text{on}}$  of RESURF SBDs can be explained by the simulation results in Fig. 2(d). The specific  $R_{\text{on}}$  of the RESURF SBDs with  $L_{AC}=30\mu\text{m}$  is calculated to be  $0.27\Omega\text{-cm}^2$ . Fig. 4(b) shows the temperature-dependent forward I-V characteristics of this RESURF SBD.

Fig. 5 benchmarks the  $R_{\text{on,sp}}$  versus  $BV$  as well as the  $BV$  versus maximum operational temperature for our devices and the state-of-the-art  $\text{Ga}_2\text{O}_3$  devices. Our device shows the highest  $BV$  and the operational temperature in multi-kilovolt  $\text{Ga}_2\text{O}_3$  devices.

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