

Crystal Anisotropy in Surface Energy Engineering (SEE) of LiTaO₃(110) Piezo-Electric for Low Temperature (< 453 K) Nano-Bonding™ to Si and α-Quartz SiO₂

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GOAL: This project’s goal is to bond *anisotropic* LiTaO₃ to Si via Nano-Bonding™¹⁻², which aims to minimize thermal stress during bonding and maximize density molecular cross-bonds via Surface Energy Engineering (SEE). SEE planarizes surfaces at three length scales to achieve direct “nano–contact”, at the nano-, micro-, and macro-scale. SEE uses Three Liquid Contact Angle Analysis (3LCAA) and the van Oss-Chaudhury-Good theory to measure hydro-affinity (HA) and surface energies (γ^T). These indicate the degree of interaction between surfaces. 3LCAA can map γ^T along different crystal orientations, thus detecting the effect of LiTaO₃’s anisotropy on SEE and Nano-Bonding™.

EXPERIMENTAL RESULTS OF NANO-BONDING™ LiTAO₃ TO SI:

Hydrophobic LiTaO₃ nano-bonds at RT in air to hydrophilic SiO₂². During nano-contacting, the nano-bonded area increases over time, but the bonding is non-uniform² as shown in Fig.1.

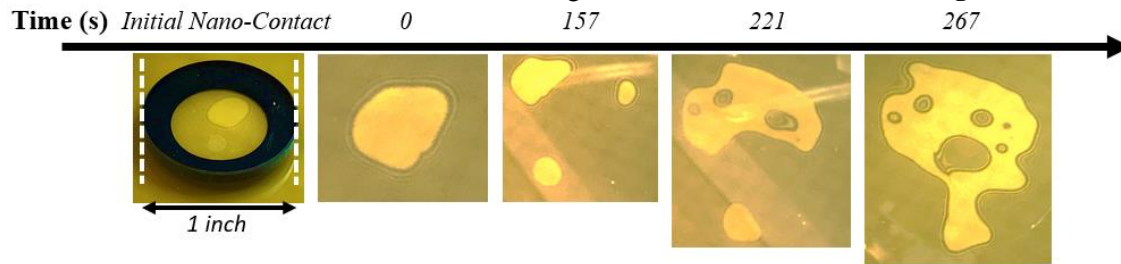


Fig. 1 Optical observation of nano-bonded area between LiTaO₃ and α-Quartz SiO₂ changing over time.

QUANTIFYING THE ROLE OF ANISOTROPY: To understand how LiTaO₃’s anisotropy plays a role in Nano-Bonding™ of LiTaO₃ to Si, elliptical distortion of the 3LCAA drops is measured in addition to water contact angle (θ_{H_2O}) and Surface Energy (γ^T) as a function of crystal orientation, as shown in Fig.2.

RESULTS AND INTERPRETATIONS:

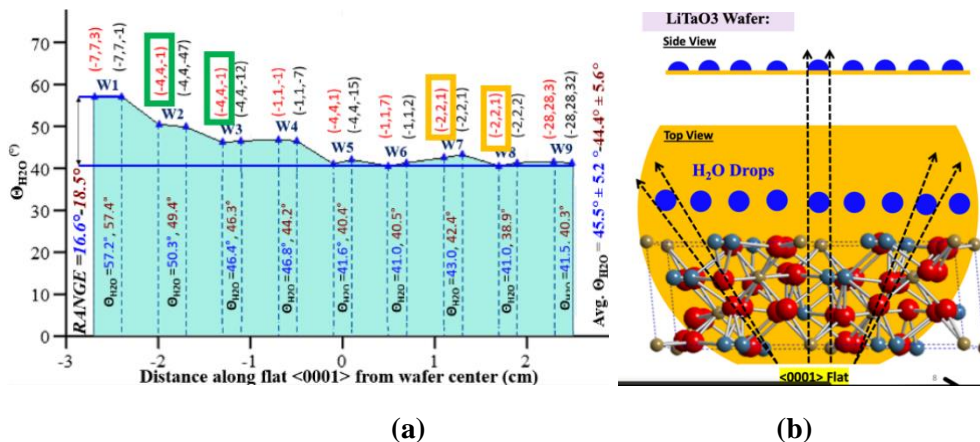


Fig.2 (a) and (b) θ_{H_2O} mapping across 15 crystal planes intersecting LiTaO₃ before SEE.

Fig. 2 (a) shows θ_{H_2O} mapping across 15 crystal planes intersecting LiTaO₃ before SEE. This shows that θ_{H_2O} varies significantly, by 40%, with crystal direction when mapped across a 150 mm (6”) LiTaO₃ (100)

¹US & International Patents Herbots N et al., n° 6,613,677 (2003); 7,851,365 (2005); 9,018,077 (2015); 9,589,801 (2017)

²Baker, B et al., J. Vac. Sci. Tech. A37, no. 4 (2019): 041101

wafer diameter. It averages $45 \pm 5.4^\circ$ along $\langle 0001 \rangle$ direction intersecting the (110) surface plane, with a range, R, of $17 \pm 1^\circ$. This is a significant variation as R is $> 3\sigma = 3 \times 5.4^\circ = 16.2^\circ$. The Miller indices of the crystal planes in (a) intersecting the location where each of the eighteen $\theta_{\text{H}_2\text{O}}$ are measured, are listed and (b) depicts how measured $\theta_{\text{H}_2\text{O}}$ correlates to crystal direction. $\theta_{\text{H}_2\text{O}}$ is measured four times for the (-4,4,-1) plane, averaging $48.1 \pm 1.8^\circ$, and the (-2,2,1) plane, averaging $41.3 \pm 1.6^\circ$. This shows that with 4 data points, the experimental error is 1.7° for these two planes, 10 x less than the measured range R.

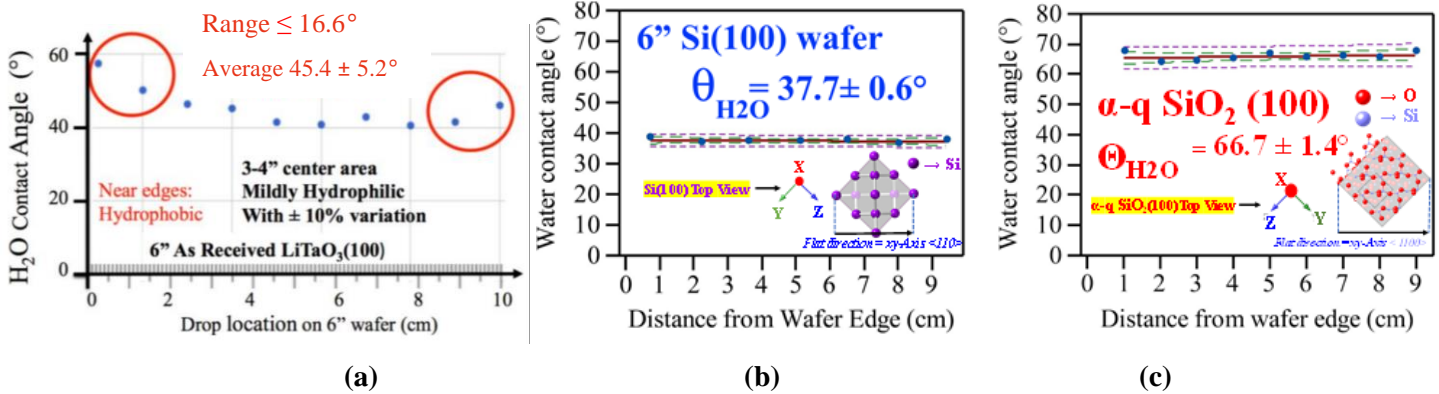


Fig. 3 Comparison of water contact angle $\theta_{\text{H}_2\text{O}}$ mapped across (a) LiTaO_3 (b) Si, and (c) α -Quartz SiO_2 wafers highlighting the uniformity of $\theta_{\text{H}_2\text{O}}$ on Si and α -Quartz SiO_2 versus the non-uniformity LiTaO_3 .

Fig.3 compares $\theta_{\text{H}_2\text{O}}$'s mapped across (a) 6" LiTaO_3 (110) along $\langle 0001 \rangle$ (b) 6" Si (100), and (c) α -Quartz SiO_2 (100). Si and SiO_2 exhibit much more uniform $\theta_{\text{H}_2\text{O}}$ mapping than LiTaO_3 by at least a full order of magnitude. LiTaO_3 's intersecting crystal planes each represents a uniquely different chemical arrangement of atoms and stoichiometric composition of elements, affecting the surface dipole and γ^T at each location on LiTaO_3 (110). In comparison, cubic Si (100) and almost cubic piezoelectric α -Quartz SiO_2 (which has also a trigonal like LiTaO_3 , but with much closer a,b,c, lattice constants) don't have very different chemical structures based off of location on wafer, thus resulting in much more $\theta_{\text{H}_2\text{O}}$ uniformity.

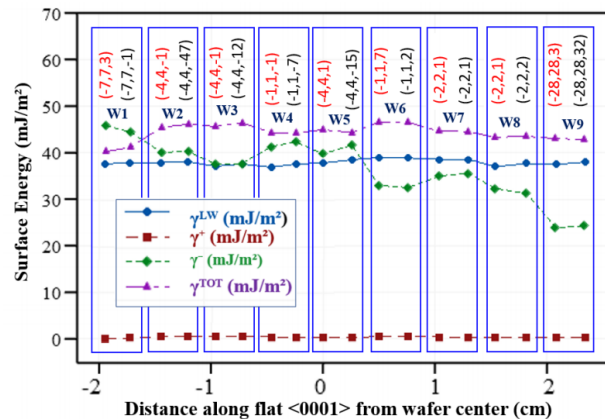


Fig. 4 Mapping of total surface energy (γ^T) and its components (γ^{LW} , γ^+ , γ^-), at each water drop position on the LiTaO_3 wafer. γ^T varies by $\pm 6\%$, which is significantly less than the 40% variation of $\theta_{\text{H}_2\text{O}}$.

Fig. 4 shows that the variations in the electron acceptor γ^- and electron donor γ^+ cancel each other out, thus resulting in a rather uniform γ^T .

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